

第5回：ニュートンの運動 方程式

必要になる基礎知識

- $\mathbf{a} = (a_x, a_y, a_z)$ の時、 $\mathbf{a} \cdot \mathbf{a}$ (= a^2 と書く)は？
- 直交座標(x, y, z)とするとき極座標(r, θ, ϕ)を用いて x, y, z を表すと？
- 直交座標(x, y, z)とするとき速度(v_x, v_y, v_z)を x, y, z で表すと？
- 極座標(r, θ, ϕ)とするとき速度(v_r, v_θ, v_ϕ)を r, θ, ϕ で表すと？

力のある場合とラグランジュの方法

力がないと？？？

空間の一様性

$$\frac{\partial \mathcal{L}(x(t), v(t))}{\partial x(t)} = 0$$

空間の等方性

$$\mathcal{L}(x(t), v(t)) = \mathcal{L}(x(t), -v(t))$$

$$\mathcal{L}(x(t), v(t)) = \frac{1}{2}m[v(t)]^2$$

力があると？？？

空間の一様性が成立しない

$$\frac{\partial \mathcal{L}(x(t), v(t))}{\partial x(t)} \neq 0$$

空間の等方性が成立しない

$$\mathcal{L}(x(t), v(t)) \neq \mathcal{L}(x(t), -v(t))$$

空間が一様でないとき

力があると？？？

空間の一様性が成立しない

$$\frac{\partial \mathcal{L}(\mathbf{x}(t), \mathbf{v}(t))}{\partial \mathbf{x}(t)} \neq 0$$

力が無いとき

$$\mathcal{L}(\mathbf{x}(t), \mathbf{v}(t)) = \frac{1}{2} m [\mathbf{v}(t)]^2$$

力があるとき

$$\mathcal{L}(\mathbf{x}(t), \mathbf{v}(t)) = \frac{1}{2} m [\mathbf{v}(t)]^2 - U(\mathbf{x}(t))$$

$$\therefore \frac{\partial \mathcal{L}(\mathbf{x}(t), \mathbf{v}(t))}{\partial \mathbf{x}(t)} = \frac{\partial \left\{ \frac{1}{2} m [\mathbf{v}(t)]^2 - U(\mathbf{x}(t)) \right\}}{\partial \mathbf{x}(t)} = -\frac{\partial U(\mathbf{x}(t))}{\partial \mathbf{x}(t)} \neq 0$$

⇒ 空間が一様でない

⇒ 力があるとき

$$\mathcal{L}(\mathbf{x}(t), \mathbf{v}(t)) = \frac{1}{2} m [\mathbf{v}(t)]^2 - U(\mathbf{x}(t))$$

ラグランジュの方法と保存力

$$\mathcal{L}(x(t), v(t)) = \frac{1}{2}m[v(t)]^2 - U(x(t))$$

ラグランジュの運動方程式

$$\frac{d}{dt}\left(\frac{\partial \mathcal{L}(x(t), v(t))}{\partial v(t)}\right) = \frac{\partial \mathcal{L}(x(t), v(t))}{\partial x(t)}$$

$$\frac{\partial \mathcal{L}(x(t), v(t))}{\partial v(t)} = \frac{\partial}{\partial v(t)}\left(\frac{1}{2}m[v(t)]^2 - U(x(t))\right) = mv(t)$$

$$\frac{\partial \mathcal{L}(x(t), v(t))}{\partial x(t)} = \frac{\partial}{\partial x(t)}\left(\frac{1}{2}m[v(t)]^2 - U(x(t))\right) = -\frac{\partial U(x(t))}{\partial x(t)}$$

$$\frac{d}{dt}(mv(t)) = -\frac{\partial U(x(t))}{\partial x(t)}$$



ニュートンの力 $\overbrace{m \frac{dv(t)}{dt}}^{\text{保存力}} = -\frac{\partial U(x(t))}{\partial x(t)}$

3次元ベクトル表記: 直交座標

$$\mathcal{L}(x(t), v(t)) = \frac{1}{2}m[v(t)]^2 - U(x(t)) \quad \left(v(t) = \frac{dx(t)}{dt} \right)$$

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$$x(t) = (x(t), y(t), z(t))$$

$$v(t) = (v_x(t), v_y(t), v_z(t))$$

$$[v(t)]^2 = v_x(t)v_x(t) + v_y(t)v_y(t) + v_z(t)v_z(t)$$

$$v(t) = \frac{dx(t)}{dt} = \left(\frac{dx(t)}{dt}, \frac{dy(t)}{dt}, \frac{dz(t)}{dt} \right)$$

$$= \frac{1}{2}m([v_x(t)]^2 + [v_y(t)]^2 + [v_z(t)]^2) - U(x(t), y(t), z(t))$$

$$= \frac{1}{2}m\left[\left(\frac{dx(t)}{dt}\right)^2 + \left(\frac{dy(t)}{dt}\right)^2 + \left(\frac{dz(t)}{dt}\right)^2\right] - U(x(t), y(t), z(t))$$

$$\mathcal{L}(x(t), v(t)) = \frac{1}{2}m\left[\left(\frac{dx(t)}{dt}\right)^2 + \left(\frac{dy(t)}{dt}\right)^2 + \left(\frac{dz(t)}{dt}\right)^2\right] - U(x(t), y(t), z(t))$$

3次元ベクトル表記: 極座標

$$\mathcal{L}(x(t), v(t)) = \frac{1}{2}m[v(t)]^2 - U(x(t)) \quad \left(v(t) = \frac{dx(t)}{dt} \right)$$

位置は… $x(t) = (r, \theta, \phi)$

速度は… $v(t) = \left(\frac{dr}{dt}, \frac{d\theta}{dt}, \frac{d\phi}{dt} \right)$

速度を計算する…

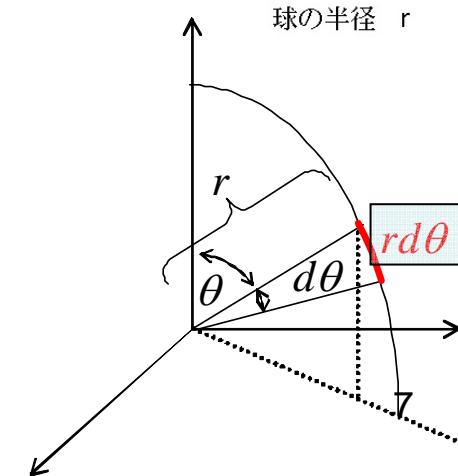
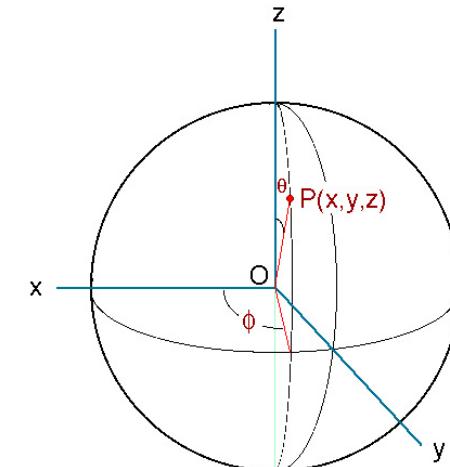
$\frac{dx}{dt}$ = **微小移動距離** / **微小経過時間** と考える

円弧の移動距離は… $rd\theta$

微小移動距離 / **微小経過時間** = $\frac{rd\theta}{dt}$



$$v_\theta = r \frac{d\theta}{dt}$$



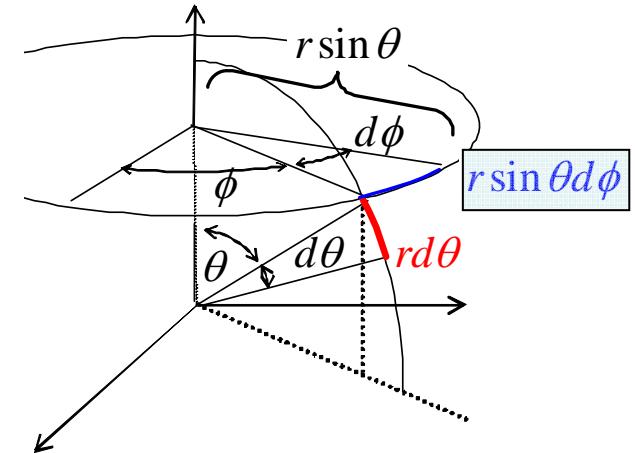
$$v(t) = \left(\frac{dr}{dt}, \frac{d\theta}{dt}, \frac{d\phi}{dt} \right) \xrightarrow{\text{}} v(t) = \left(\frac{dr}{dt}, r \frac{d\theta}{dt}, \frac{d\phi}{dt} \right)$$

もうひとつの円弧の移動距離は…

$$r \sin \theta d\phi$$

$$\frac{\text{微小移動距離}}{\text{微小経過時間}} = \frac{r \sin \theta d\phi}{dt}$$

$$v_\phi = r \sin \theta \frac{d\phi}{dt}$$



$$v(t) = \left(\frac{dr(t)}{dt}, r \frac{d\theta(t)}{dt}, r \sin \theta \frac{d\phi(t)}{dt} \right)$$

$$\boxed{\boldsymbol{v}(t) = \left(\frac{dr(t)}{dt}, r(t) \frac{d\theta(t)}{dt}, r(t) \sin \theta(t) \frac{d\phi(t)}{dt} \right)}$$

$$\begin{aligned}
\mathcal{L}(\boldsymbol{x}(t), \boldsymbol{v}(t)) &= \frac{1}{2}m \left[\boldsymbol{v}(t) \right]^2 - U(\boldsymbol{x}(t)) \left(\boldsymbol{v}(t) = \frac{d\boldsymbol{x}(t)}{dt} \right) \\
\boldsymbol{x}(t) &= (r(t), \theta(t), \phi(t)) \\
\boldsymbol{v}(t) &= (\boldsymbol{v}_r(t), \boldsymbol{v}_\theta(t), \boldsymbol{v}_\phi(t)) \\
\boldsymbol{v}(t) &= \frac{d\boldsymbol{x}(t)}{dt} = \left(\frac{dr}{dt}, r \frac{d\theta}{dt}, r \sin \theta \frac{d\phi}{dt} \right) \\
&= \frac{1}{2}m \left([\boldsymbol{v}_x(t)]^2 + [\boldsymbol{v}_y(t)]^2 + [\boldsymbol{v}_z(t)]^2 \right) - U(r(t), \theta(t), \phi(t)) \\
&= \frac{1}{2}m \left[\left(\frac{dr(t)}{dt} \right)^2 + \left(r \frac{d\theta(t)}{dt} \right)^2 + \left(r \sin \theta \frac{d\phi(t)}{dt} \right)^2 \right] - U(r(t), \theta(t), \phi(t))
\end{aligned}$$

$$\boxed{\mathcal{L}(\boldsymbol{x}(t), \boldsymbol{v}(t)) = \frac{1}{2}m \left[\left(\frac{dr(t)}{dt} \right)^2 + \left(r(t) \frac{d\theta(t)}{dt} \right)^2 + \left(r(t) \sin \theta(t) \frac{d\phi(t)}{dt} \right)^2 \right] - U(r(t), \theta(t), \phi(t))}$$

3次元ベクトル表記(まとめ)

$$\mathcal{L}(x(t), v(t)) = \frac{1}{2}m [v(t)]^2 - U(x(t)) \quad \left(v(t) = \frac{dx(t)}{dt} \right)$$

$$\mathcal{L}(x(t), v(t)) = \frac{1}{2}m \left[\left(\frac{dx(t)}{dt} \right)^2 + \left(\frac{dy(t)}{dt} \right)^2 + \left(\frac{dz(t)}{dt} \right)^2 \right] - U(x(t), y(t), z(t))$$

$$x(t) = (x(t), y(t), z(t))$$

$$v(t) = \frac{dx(t)}{dt} = \left(\frac{dx(t)}{dt}, \frac{dy(t)}{dt}, \frac{dz(t)}{dt} \right)$$

$$\mathcal{L}(x(t), v(t)) = \frac{1}{2}m \left[\left(\frac{dr(t)}{dt} \right)^2 + \left(r(t) \frac{d\theta(t)}{dt} \right)^2 + \left(r(t) \sin \theta(t) \frac{d\phi(t)}{dt} \right)^2 \right] - U(r(t), \theta(t), \phi(t))$$

$$x(t) = (r(t), \theta(t), \phi(t))$$

$$v(t) = \frac{dx(t)}{dt} = \left(\frac{dr}{dt}, r \frac{d\theta}{dt}, r \sin \theta \frac{d\phi}{dt} \right)$$